

# **COMBUSTION FUNDAMENTALS**

## **Premixed and Non-Premixed Turbulent Flames**

# EXAMPLE: SPARK IGNITION ENGINES

- Even though fuel is introduced as a liquid, spark-ignition engines fuels are highly volatile, and liquid has time to vaporize and thoroughly mix with air before mixture ignited by spark
- Combustion duration is an important parameter in operation of spark-ignition engines and is controlled by **turbulent flame speed** and distribution of combustion volume
- Compact combustion chambers produce short combustion durations
- Combustion duration governs lean-limit of stable operation, tolerance to exhaust gas recirculation, thermal efficiency, and production of  $\text{NO}_x$  emissions

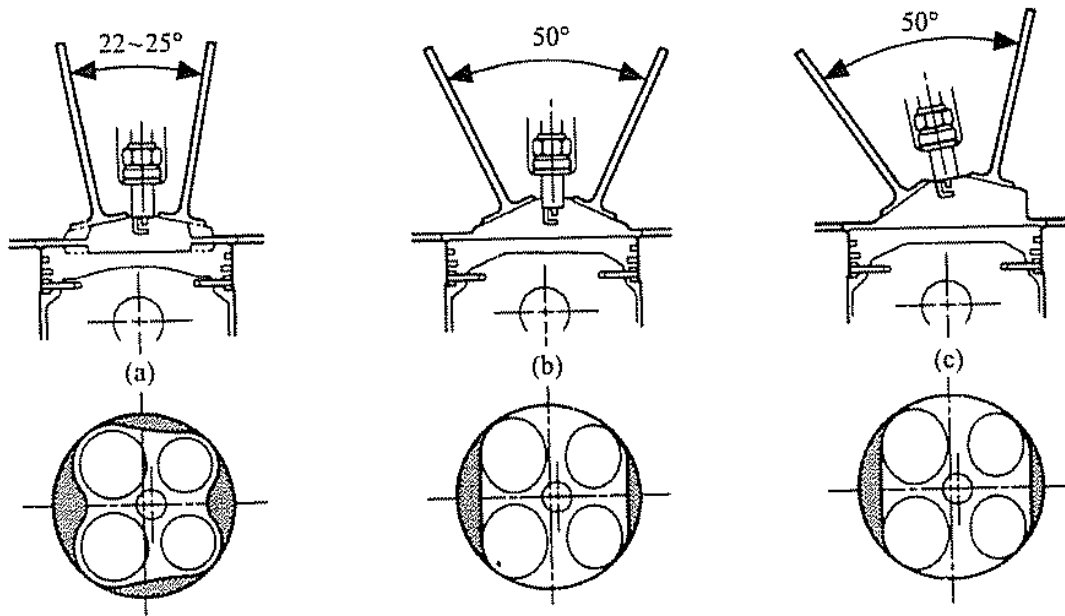
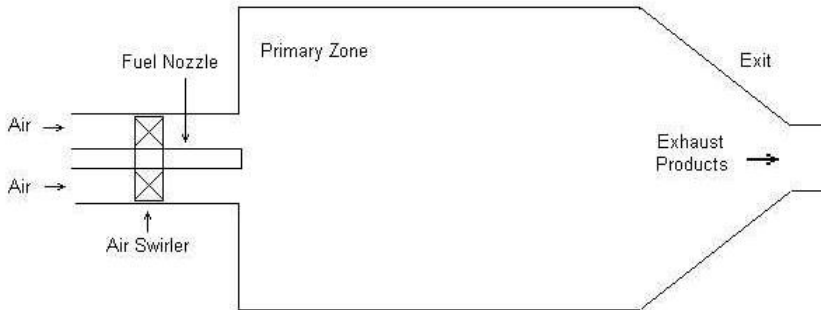


Figure 12.1 Various configurations of four-valve, spark-ignition engine combustion chambers  
| SOURCE: Courtesy of General Motors Corporation

# EXAMPLE: GAS TURBINE ENGINES

- Engines are being more and more used for ground based power
- Current combustor design is largely influenced by the need to control soot, CO, and NO<sub>x</sub>
- Older engines employed purely non-premixed (diffusion) combustors
  - Near stoichiometric burning primary zone
  - Secondary air to complete combustion and reduce temperature prior to entering turbine
- Some current designs use some premixing to avoid high temperature, NO<sub>x</sub> formation zones
  - However, there are drawbacks with this design:
    - Flame stability
    - CO emissions
    - Ratio of maximum to minimum flow rates (called turndown ratio)



LPP Combustor

LPP: Lean premixed prevaporized

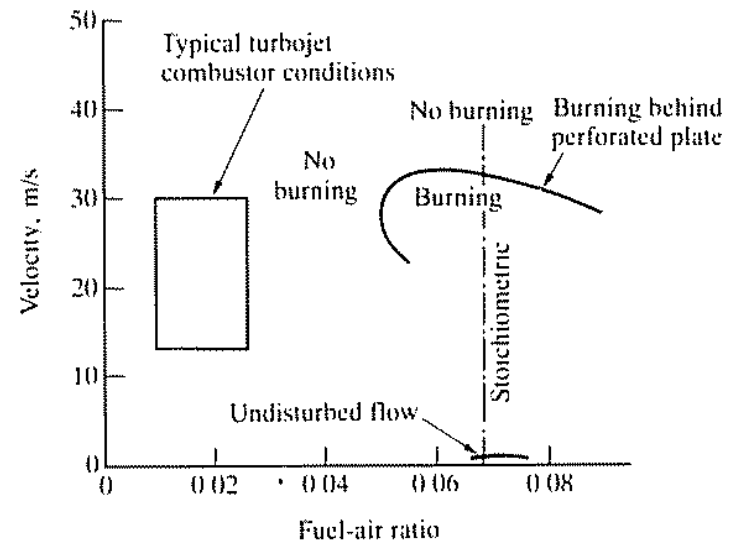
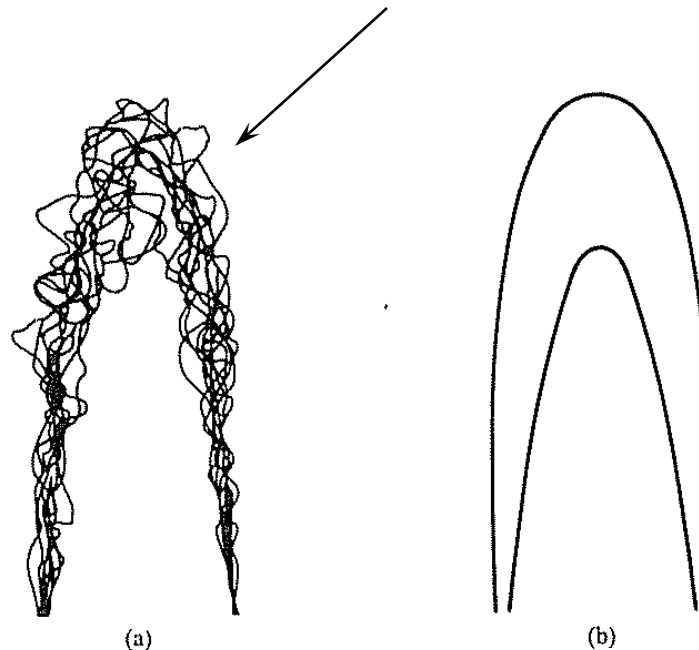


FIGURE 6.24 Dependence of flame speed on fuel-air ratio (Adapted from Olson, Childs, and Jonash [14])

# STRUCTURE OF TURBULENT PREMIXED FLAMES

- Instantaneous superimposed contours of convoluted thin reaction zones
  - Obtained using schlieren photography at different instants in time
  - Large folds near top of flame
  - Position of reaction zone moves rapidly in space, producing a time-averaged view that gives appearance of a thick reaction zone, which is called turbulent flow brush
  - Instantaneous view shows that actual reaction front is relatively thin, as in laminar premixed flame
    - Sometimes referred to as laminar flamelets



**Figure 12.6** (a) Superposition of instantaneous reaction fronts obtained at different times  
(b) Turbulent flame "brush" associated with a time-averaged view of the same flame  
| SOURCE: (a) After Ref. [15] from Ref [19], reprinted by permission of Academic Press

# DEFINITION OF TURBULENT FLAME SPEED, $S_t$

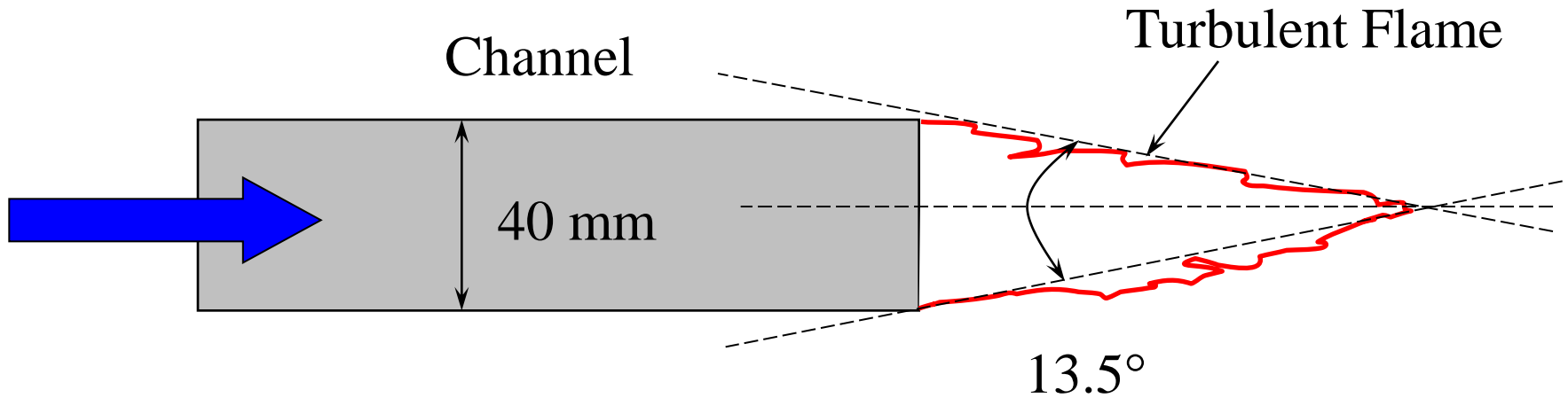
- Recall that **laminar flames** have a propagation velocity,  $S_L$ , that depends **uniquely** on **thermal** and **chemical properties of the mixture**
- **Turbulent flame** flames have a propagation velocity that depends on the **character of flow**, as well as on **mixture properties**
- For an observer traveling with the flame, we can define a turbulent flame speed,  $S_t$ , as the velocity at which unburned mixture enters the flame zone in a direction normal to the flame
  - Flame surface is represented as some time mean quantity
  - Instantaneous portions of the high temperature reaction zone may be largely fluctuating
  - Usually determined from measurements of reactant flowrates
- Turbulent flame speed can be expressed as:

$$S_t = \frac{\dot{m}}{A\rho_u}$$

- Experimental determinations of turbulent flame speeds are complicated by determining a suitable flame area, for what are usually thick and frequently curved flames
- This ambiguity results in considerable uncertainty in measurements of turbulent flame velocities

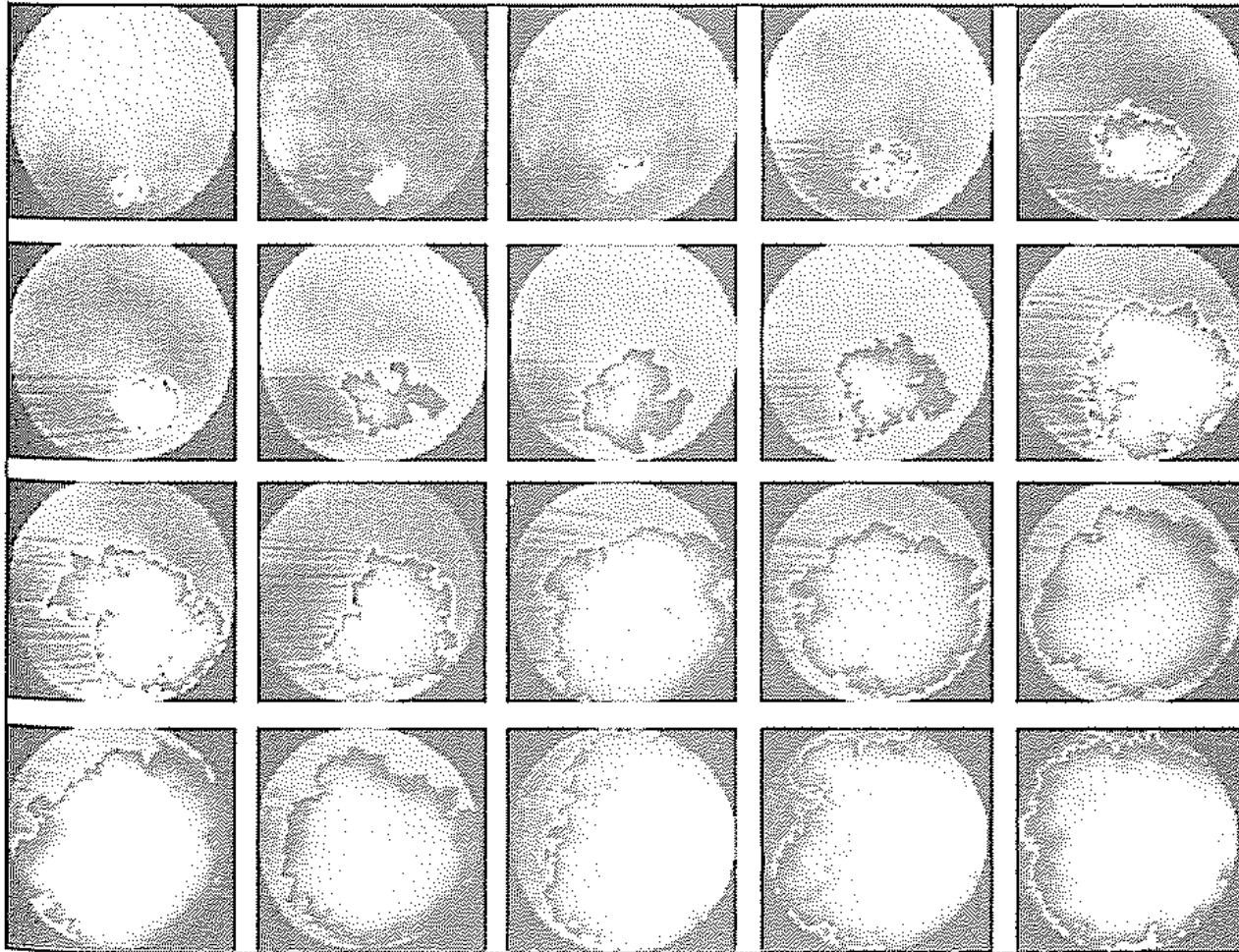
## EXAMPLE: FROM EXPERIMENT

- An air-fuel mixture passes through a 40 mm by 40 mm flow channel with a flame anchored at channel exit along top and bottom walls, as shown below



- Quartz side walls contain flame beyond exit, while top and bottom are open, so assume flame forms a wedge shape
  - Mean flow velocity is 70 m/s
  - Density of unburned gas is 1.2 kg/m<sup>3</sup>
  - Wedge shaped flame has an angle of 13.5°, which was estimated from time averaged photographs
  - MW = 29
- Estimate turbulent burning velocity at this condition

# EXAMPLE: SPARK IGNITION ENGINE VIEW



- Visualization of turbulent flame propagation in a spark-ignition engine operating at 1,200 RPM
- Images represent a planar slice through the combustion chamber with sequence starting soon after ignition (upper left photo) and proceeding until flame comes to cylinder walls
- The flame structure in these photos is in the ‘wrinkled laminar flame regime’
- Speeding up the engine to 2,400 RPM would produce a flame with ‘pockets’ or ‘islands’ of burned and unburned gases, which is given the structural name ‘Flamelets in eddies regime’

# 3 FLAME REGIMES

- Various length scales exist simultaneously in a turbulent flow
  - Smallest is called the Kolmogorov microscale,  $l_K$ , which represents smallest eddies in flow
    - Eddies rotate rapidly and have high vorticity ( $\omega = \text{DEL} \times V$ ), which results in dissipation of fluid kinetic energy into internal energy (fluid friction results in a temperature rise of fluid)
  - Integral scale,  $l_0$ , characterizes largest eddies
- Basic structure of turbulent flame governed by relationships of  $l_K$  and  $l_0$  to laminar flame thickness,  $\delta_L$
- Laminar flame thickness characterizes thickness of reaction zone controlled by molecular (not turbulent) transport of heat and mass

## 1. **Wrinkled laminar flame regime: $\delta_L \leq l_K$**

- When the flame thickness is much thinner than the smallest scale of turbulence, the turbulent motion can only wrinkle or distort the thin laminar flame zone
- Criterion for existence of a wrinkled laminar flame is referred to as [Williams-Klimov criterion](#)

## 2. **Distributed reaction regime: $\delta_L > l_0$**

- If all scales of turbulent motion are smaller than reaction zone thickness, transport within reaction zone is no longer governed solely by molecular processes, but also by turbulence
- Criterion for existence of a distributed reaction zone is called [Damköhler criterion](#)

## 3. **Flamelets-in-eddies regime: $l_0 > \delta_L > l_K$**



## DAMKÖHLER NUMBER, $Da$

- Important dimensionless number in combustion,  $Da$
- Represents a ratio of characteristic flow time to characteristic chemical time =  $\tau_{\text{flow}}/\tau_{\text{chem}}$
- In premixed flames, the following time scales are particularly useful
  - Flow time,  $\tau_{\text{flow}} \equiv l_0/v'_{\text{RMS}}$
  - Chemical time based on a laminar flame,  $\tau_{\text{chem}} \equiv \delta_L/S_L$

$$Da = \frac{\left( \frac{l_0}{v'_{\text{RMS}}} \right)}{\left( \frac{\delta_L}{S_L} \right)} = \left( \frac{l_0}{\delta_L} \right) \left( \frac{S_L}{v'_{\text{RMS}}} \right)$$

- IF  $Da \gg 1$  reaction rates are very fast in comparison with fluid mixing rates
  - Called fast chemistry regime
- IF  $Da \ll 1$  reaction rates are slow in comparison with mixing rates
- Note if fix length scale ratio,  $Da$  falls as turbulence intensity goes up

# GOVERNING NON-DIMENSIONAL NUMBERS

$$\frac{l_K}{\delta_L}$$

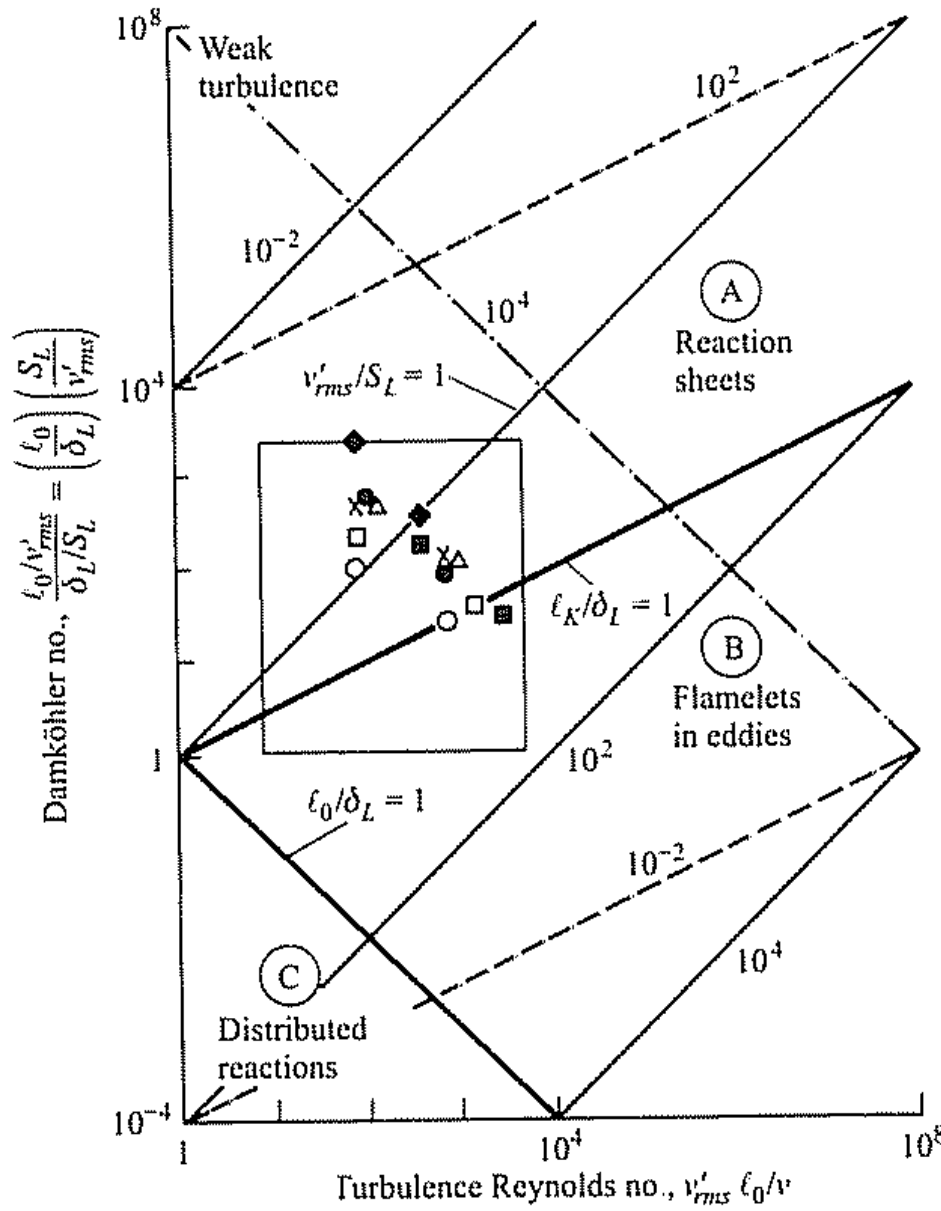
$$\frac{l_0}{\delta_L}$$

$$\text{Re}_{l_0} \equiv \frac{v'_{RMS} l_0}{\nu}$$

$$Da = \frac{l_0 / v'_{RMS}}{\delta_L / S_L} = \left( \frac{l_0}{\delta_L} \right) \left( \frac{S_L}{v'_{RMS}} \right)$$

$$\frac{v'_{RMS}}{S_L}$$

# IMPORTANT PARAMETERS CHARACTERIZING TURBULENT PREMIXED COMBUSTION



- What flame regime do practical devices fall under?
- Conditions satisfying Williams-Klimov criterion for wrinkled flames lie above solid line ( $l_K = \delta_L$ )
- Conditions satisfying Damköhler criterion for distributed reactions fall below solid line ( $l_0 = \delta_L$ )
- Thin reaction sheets can only occur for  $Da > 1$ , depending on  $Re$ , which indicates that regime is characterized by fast chemistry as compared with fluid mixing
- Box shows spark ignition engine data

# COMMENTS ON WRINKLED LAMINAR FLAME REGIME

- Chemical reactions occur in thin sheets,  $Da > 1$ , fast chemistry region
- Only effect of turbulence is to wrinkle flame, resulting in an increased flame area

$$\frac{S_t}{S_L} = 1 + \frac{v'_{RMS}}{S_L} \quad \text{Damköhler}$$

$$\frac{S_t}{S_L} = 3.5 \left( \frac{v'_{RMS}}{S_L} \right)^{0.7} \quad \text{Klimov}$$

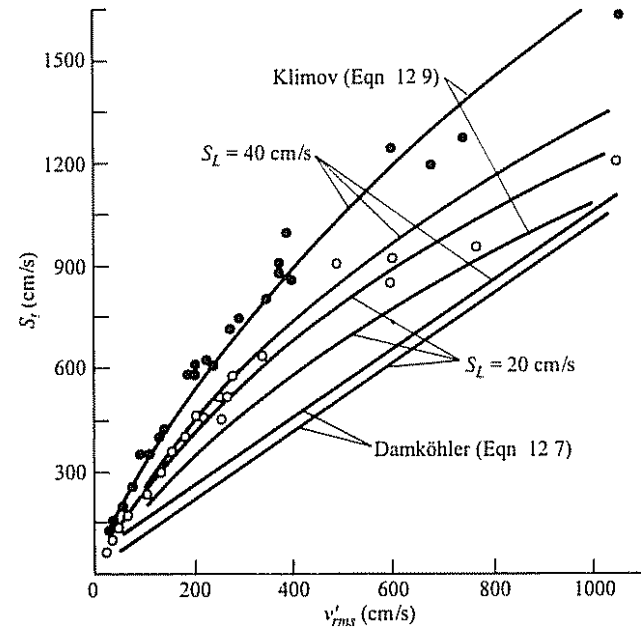


Figure 12.10 Experimental data for  $S_t$  versus  $v'_{rms}$ , compared with wrinkled laminar-flame theories of turbulent flame propagation  
 | SOURCE: Data from Ref [23]

- Example: Laser anemometry is used to measure the mean and fluctuating velocities in a spark ignition engine. Estimate the turbulent flame speed for  $v'_{RMS} = 3 \text{ m/s}$ ,  $P = 5 \text{ atm}$ ,  $T_u = 500 \text{ }^\circ\text{C}$ ,  $\phi = 1.0$  for a propane-air mixture, and the mass fraction of the residual burned gases mixed with fresh air is 0.09.

# FLAME SPEED CORRELATIONS FOR SELECTED FUELS

- One of most useful correlations for laminar flame speed,  $S_L$ , given by Metghalchi and Keck
  - Determined experimentally over a range of temperatures and pressures typical of those found in reciprocating IC engines and gas-turbine combustors

$$S_L = S_{L,ref} \left( \frac{T_u}{T_{u,ref}} \right)^\gamma \left( \frac{P}{P_{ref}} \right)^\beta (1 - 2.1Y_{dil})$$

$$S_{L,ref} = B_M + B_2(\phi - \phi_M)^2$$

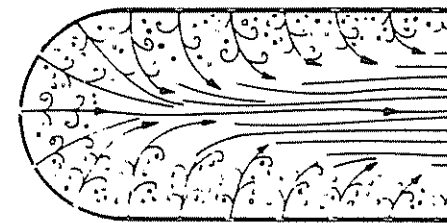
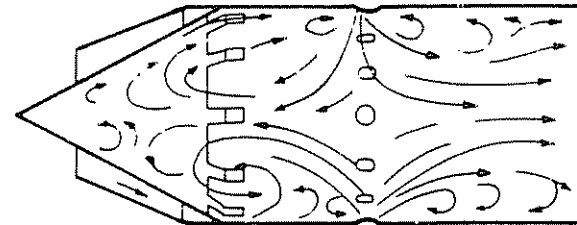
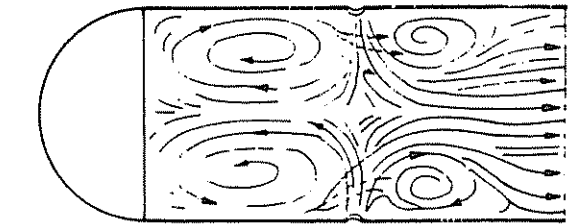
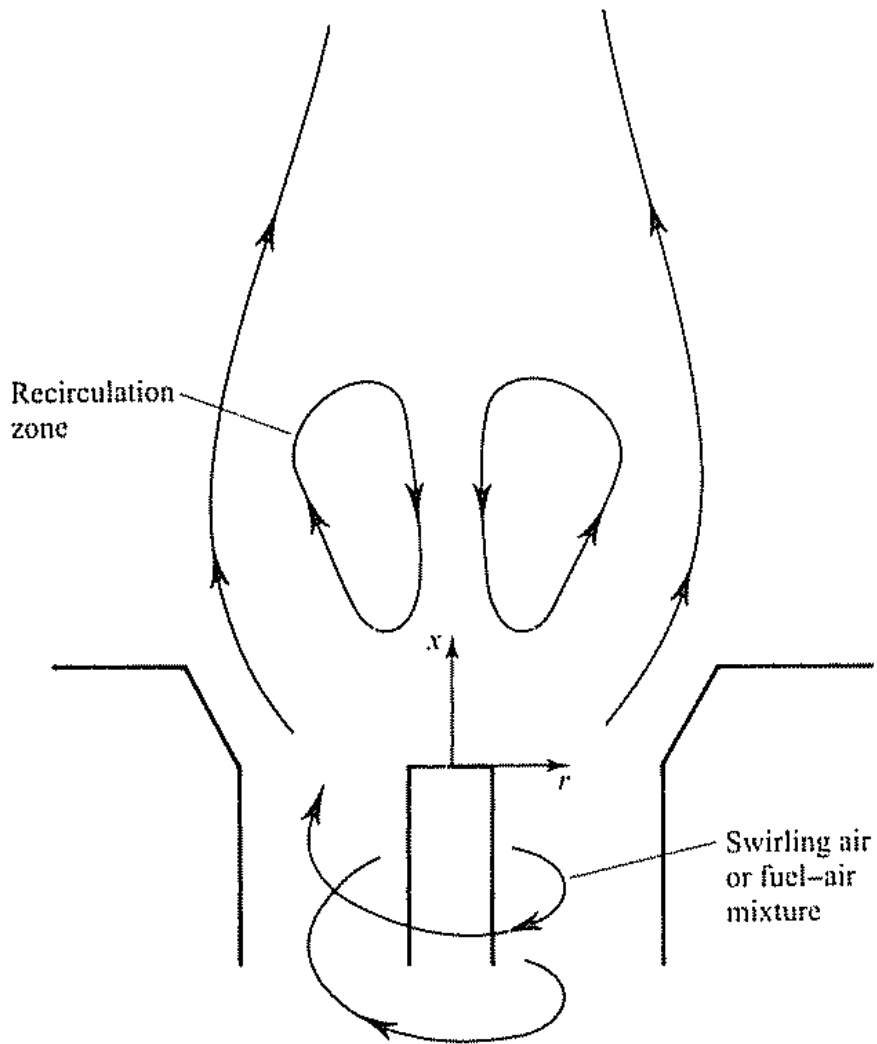
$$\gamma = 2.18 - 0.8(\phi - 1)$$

$$\beta = -0.16 + 0.22(\phi - 1)$$

Fuel	$\Phi_M$	$B_M$ (cm/s)	$B_2$ (cm/s)
Methanol	1.11	36.92	-140.51
Propane	1.08	34.22	-138.65
Isooctane	1.13	26.32	-84.72
RMFD-303	1.13	27.58	-78.34

- **EXAMPLE:** Employ correlation of Metghalchi and Keck to compare laminar flame speed gasoline (RMFD-303)-air mixtures with  $\phi = 0.8$  for 3 cases:
  1. At reference conditions of  $T = 298$  K and  $P = 1$  atm
  2. At conditions typical of a spark ignition engine operating at  $T = 685$  K and  $P = 18.38$  atm
  3. At same conditions as (2) but with 15 percent (by mass) exhaust gas recirculation

# INFLUENCE OF SWIRL



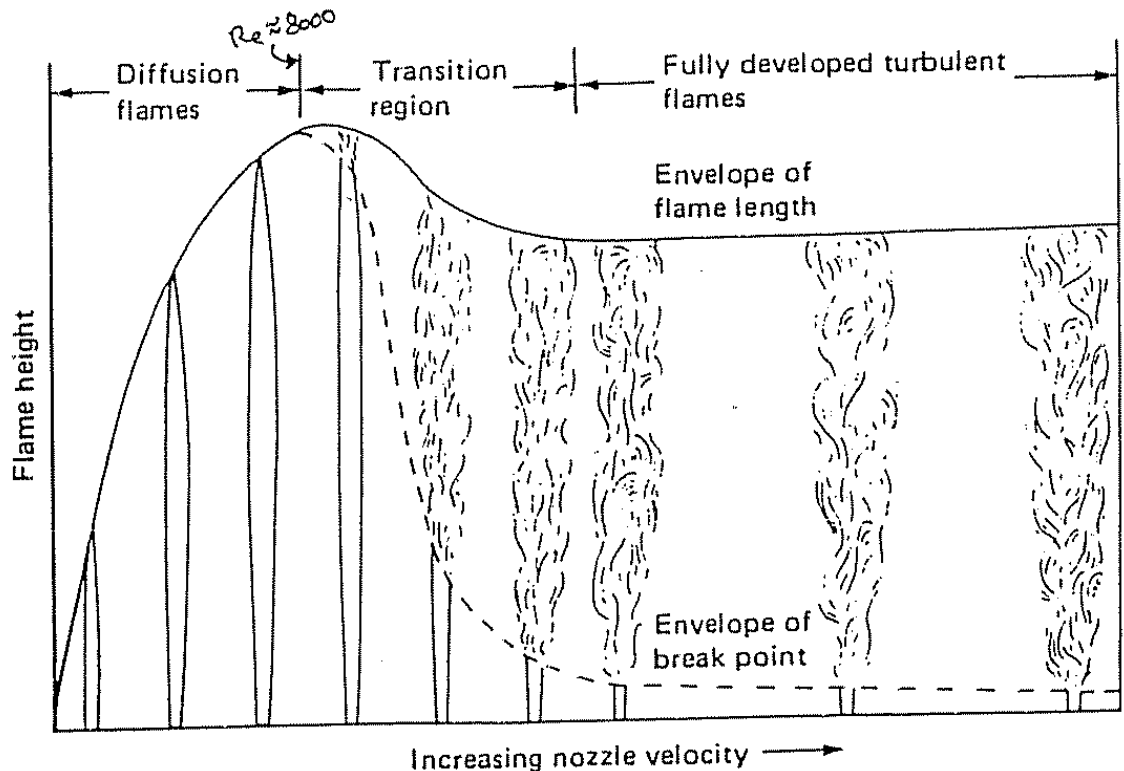
**Figure 12.17** Flow patterns for can combustors with (a) single row of holes with enclosed end, (b) shrouded cone, and (c) multiple rows of holes

[ SOURCE: From Ref [31], reprinted by permission of The Combustion Institute

# OVERVIEW: TURBULENT NON-PREMIXED (DIFFUSION) FLAMES

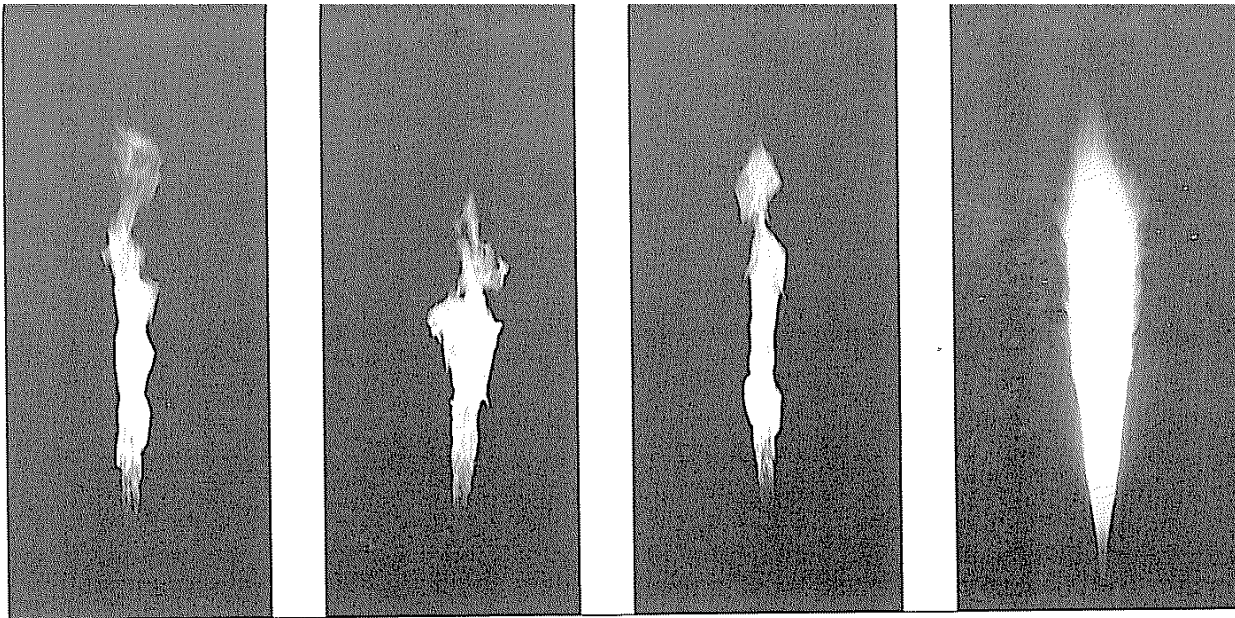
- Turbulent non-premixed flames are employed in most practical devices as they are easier to control
- With pollutants a major concern, this advantage can become a liability
  - Less ability to control pollutant formation or ‘tailor’ flow field
- Examples
  - For low NO<sub>x</sub> in a gas turbine combustor usually new trend is to use premixed primary zones
  - Flames stabilized behind bluff bodies in afterburners for military aircraft
  - Liquid fuel sprays in diesel engines

- Engineering challenges
  - Flame shape and size
  - Flame holding and stability
  - Heat transfer
  - Pollutant emissions



# COMMENTS ON JET FLAMES

- Turbulent non-premixed flames also have wrinkled, contorted and brushy looking edges, just like premixed flames
- Non-Premixed flames are usually more luminous than premixed flames due to soot within the flame
- No universal definition of flame length
  - Averaging of individual flame lengths from photographs
  - Measuring location of average peak centerline temperature using thermocouples
  - Measuring location where mean mixture fraction on axis is stoichiometric using gas sampling
  - In general, visible flame lengths tend to be larger than those based on temperature or concentration measurements

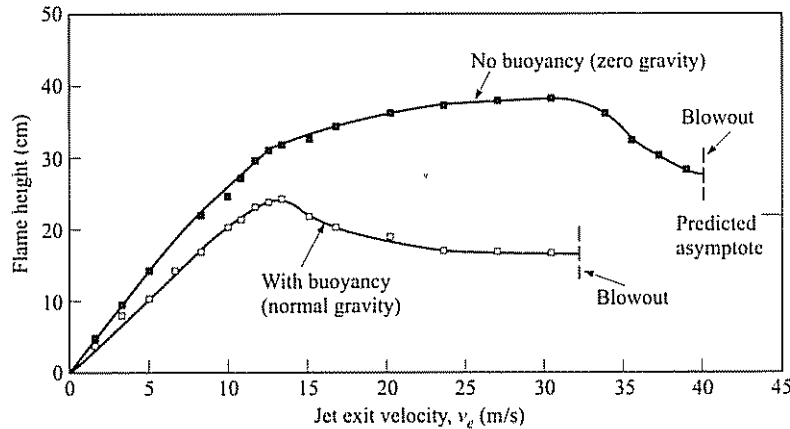


**Figure 13.3** Instantaneous and time-averaged photographs of turbulent jet diffusion flames ( $C_2H_4$ -air,  $d_j = 2.18$  mm). Because of the low luminosity at the base of the flame, this region is not visible on the instantaneous images (the first three photographs).



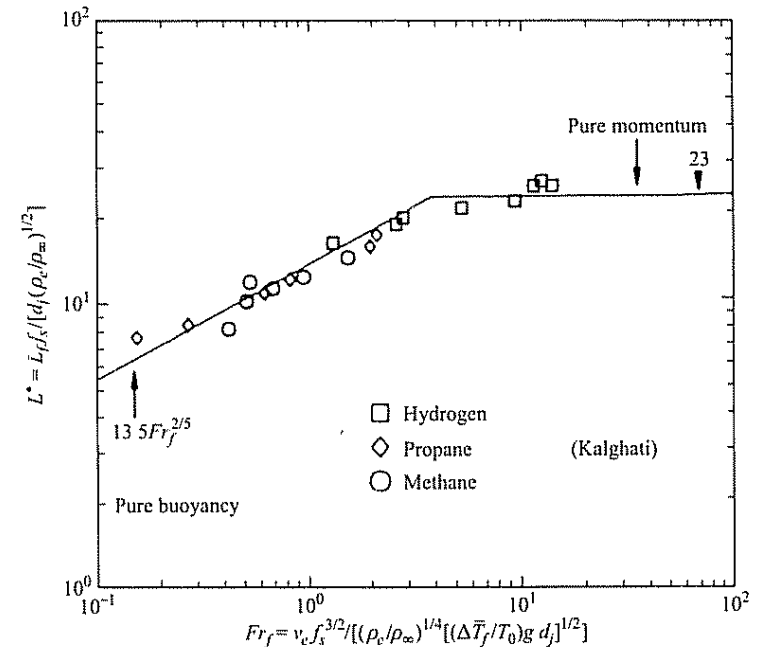
# FACTORS THAT AFFECT FLAME LENGTH, $L_f$

- Factors affecting flame length (vertical flames issuing into a still environment)
  - Relative importance of initial jet momentum flux and buoyant forces acting on flame,  $Fr$ 
    - Recall Froude number,  $Fr$ , was used to establish momentum controlled vs. buoyancy controlled flow regimes for laminar jet flames
    - $Fr \gg 1$ : flames are dominated by initial jet momentum, which controls mixing and velocity field within flame
    - $Fr \ll 1$ : flames are dominated by buoyancy
  - Stoichiometric mixture fraction,  $f_s = 1/((A/F)_s + 1)$
  - Ratio of nozzle fluid to ambient gas density,  $\rho_e/\rho_\infty$
  - Initial jet diameter,  $d_j$



**Figure 13.11** Comparison of jet flame heights with and without buoyancy ( $C_3H_8$ -air,  $d_j = 0.8$  mm) Nonbuoyant conditions result from tests being conducted with near-zero ( $< 10^{-5}$  g) gravitational acceleration. Flames are nominally laminar at exit velocities up to approximately 10 m/s.

| SOURCE: Data from Bahadori *et al* [19]



**Figure 13.12** Flame lengths for jet flames correlated with flame Froude number  
 | SOURCE: Reprinted by permission of Elsevier Science, Inc., from Ref. [15]. © 1993, The Combustion Institute

# USEFUL CORRELATIONS AND EXAMPLE

$$Fr = \frac{v_e f_s^{1.5}}{\left(\frac{\rho_e}{\rho_\infty}\right)^{0.25} \left[\frac{\Delta T_f}{T_\infty} g d_j\right]^{0.5}}$$

Useful definition of Fr

$\Delta T_f$  = temperature rise from combustion

$$d_j^* = d_j \left(\frac{\rho_e}{\rho_\infty}\right)^{0.5}$$

Combination of density ratio and jet diameter

Called momentum diameter

$$L^* \equiv \frac{L_f f_s}{d_j \left(\frac{\rho_e}{\rho_\infty}\right)^{0.5}} = \frac{L_f f_s}{d_j^*}$$

Dimensionless flame length,  $L^*$

From correlated data (on previous slide)

$$L^* = \frac{13.5 Fr^{2/5}}{(1 + 0.07 Fr^2)^{1/5}}$$

Buoyancy dominated regime,  $Fr < 5$

$$L^* = 23$$

Momentum dominated regime

- Simple Example: Estimate flame length for a propane jet flame in air at ambient conditions. Propane mass flow rate is  $3.7 \times 10^{-3}$  kg/s and nozzle exit diameter is 6 mm. Propane density is  $1.85$  kg/m<sup>3</sup>

# LIFTOFF AND BLOWOUT

Kalghatgi correlation to estimate blowout flow rate for jet flames

$$\frac{v_e}{S_{L,\max}} \left( \frac{\rho_e}{\rho_\infty} \right)^{1.5} = 0.017 \text{Re}_H \left( 1 - 3.5 \times 10^{-6} \text{Re}_H \right)$$

$$\text{Re}_H = \frac{\rho_e S_{L,\max} H}{\mu}$$

$$H = 4 \left[ \frac{Y_{F,e}}{Y_{F,\text{stoic}}} \left( \frac{\rho_e}{\rho_\infty} \right)^{0.5} - 5.8 \right] d_j$$

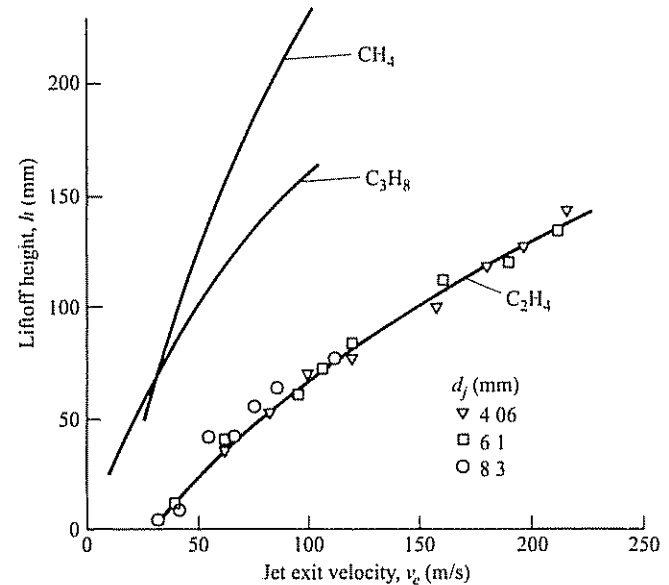


Figure 13.16 Liftoff height versus jet exit velocity for methane, propane, and ethylene jet flames.  
| SOURCE: After Kalghatgi [34]

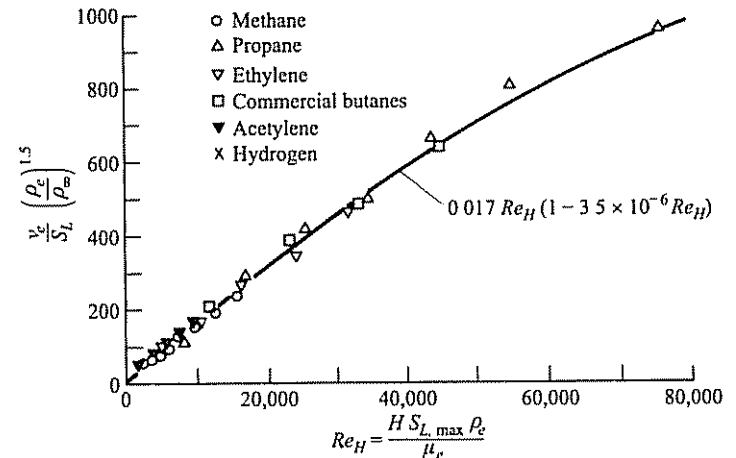


Figure 13.17 Universal blowout stability curve from Kalghatgi [35].  
| SOURCE: Reprinted with permission © 1981, Gordon & Breach Publishers

## SIMPLE EXAMPLES CONTINUED

- From previous: Estimate flame length for a propane jet flame in air at ambient conditions.
  - Propane mass flow rate is  $3.7 \times 10^{-3}$  kg/s and nozzle exit diameter is 6 mm
  - Propane density at nozzle exit is  $1.85$  kg/m<sup>3</sup>
- For same heat release rate and nozzle exit diameter, determine flame length when fuel is methane and compare with propane flame length
  - Density of methane is  $0.6565$  kg/m<sup>3</sup>
- For propane jet flame, determine blowoff velocity and estimate liftoff height at incipient blowoff condition
  - Viscosity of propane is  $8.26 \times 10^{-6}$  N s/m<sup>2</sup>.
  - To estimate liftoff height, use figure 13.16 on previous slide